

Structural and electrical properties of La-Sr-Mn-O ceramics with Bi³⁺ content for thermistor devices

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La_{0.7}Sr_{0.3-x}Bi_xMnO₃ (LSBMO) (0.025 ≤ x ≤ 0.15) films are fabricated by a sol-gel method and spin-coating method. All of the fabricated films have typical polycrystalline perovskite structures. In addition, no preferential orientation or impurity could be observed. The surface SEM images show almost homogeneous and uniform microstructures and the average thickness for the LSBMO films is approximately 230 to 260 nm. When the resistivity is measured at room temperature, the TCR and B_{85/25}-values gradually increases as the amount of Bi³⁺ increased, and the highest values of 0.023 mΩ·cm, 0.31%/°C and 299 K are obtained in the La_{0.7}Sr_{0.225}Bi_{0.075}MnO₃ thin films. Polycrystalline LSBMO films follow the variable-range hopping (VRH) conduction model as they show electron-electron scattering and electron-phonon interactions for thermal stimulation.

Keywords : La_{0.7}Sr_{0.3-x}Bi_xMnO₃ films, Sol-gel method, Structural properties, Electrical properties, Thermistor.

Introduction

Perovskite-structured manganese-based R_{1-x}A_xMnO₃ (R=trivalent rare-earth, A=divalent alkaline-earth ions) compounds exhibit colossal magnetoresistance (CMR) caused by the transition between a low-temperature ferromagnetic metallic and a high-temperature paramagnetic insulation state depending on constituent elements, applied temperature and magnetic fields [1]. These properties result from Mn³⁺ with 3d⁴ (t_{2g}³, e_g¹) electronic arrays converting to Mn⁴⁺ with 3d³ (t_{2g}³, e_g⁰) electron arrays to maintain electrical neutrality because A²⁺ ions are generally substituted at R³⁺ sites. High electrical conduction is exhibited by the double exchange (DE) of electrons with strong on-site Hund's coupling between the partially filled d-shell of Mn³⁺ and Mn⁴⁺ in the perovskite MnO₆ oxygen octahedral structure [2]. These CMR characteristics have given rise to a variety of studies focused on new functional oxide material devices, such as magnetic information storage devices, sensors and spintronics devices [3, 4]. One of the best conductive oxide materials is La_{1-x}Sr_xMnO₃, which has an electrical conductivity that is greatly affected by the Jahn-Teller distortion of the

unit lattice due to a difference in the ion radius of Mn⁴⁺/Mn³⁺, and La³⁺ (0.136 nm) and Sr²⁺ (0.144 nm) according to the composition ratio of Sr²⁺ [5, 6]. Furthermore, the electrical conduction characteristics depend heavily on microstructural characteristics such as the grain of the grains, grain boundaries, pores and oxygen attackers in accordance with the manufacturing of bulk or thin films [7].

La_{0.7}Sr_{0.3}MnO₃ (LSMO) exhibited the excellent electrical conductivity and CMR properties at room temperature due to its semi-metallic ferromagnetic characteristics [8]. In this study, we fabricated La_{0.7}Sr_{0.3-x}Bi_xMnO₃ by substituting a portion of Sr²⁺ ions with Bi³⁺ ions, which have different atomic and ionic radii, using the sol-gel method. We then observed the structural and electrical properties as a function of composition ratio to investigate the possibility of application as a thermistor devices or bolometer-type infrared detectors operating at room temperature.

Experimental

Manganese acetate (Mn(CH₃COO)₂·4H₂O), lanthanum acetate (La(CH₃COO)₃·xH₂O), strontium acetate (Sr(CH₃COO)₂) and bismuth acetate (Bi(CH₃COO)₃) were dissolved into acetic acid, ethyl alcohol and distilled water co-solvents to drive the La_{0.7}Sr_{0.3-x}Bi_xMnO₃ (0.025 ≤ x ≤ 0.15) (LSBMO) precursor solutions. The precursor solutions were aged for 24 hours before deposition.

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LSBMO coating solutions were spin-coated on cleaned Pt/Ti/SiO₂/Si substrates at 4,000 rpm for 30 seconds. The coated thin films were dried and annealed at 200 °C for 5 min. and at 400 °C for 10 min., respectively. The spin-coating, drying and pyrolysis cycle was repeated six times. Finally, the films were crystallized at 800 °C for 60 min. in an oxygen atmosphere. The crystallographic and structural properties of the films were measured using X-ray diffraction (XRD), field-emission scanning electron microscopy (FE-SEM), and X-ray photoelectron spectroscopy (XPS), and the electrical properties were measured by the Van der Pauw method using an electrometer (Keithley 6517A, USA).

Results and Discussion

Fig. 1 shows the X-ray diffraction patterns according to the amount of Bi³⁺ in the LSBMO films. All of the films have a typical perovskite structure with no preferential orientations or impurities. Perovskite LaMnO₃ adopts an orthorhombic structure at room temperature, and shows a rhombohedral structure by distortion of the ABO₃ unit cell as Sr²⁺ ions with a larger ionic radius is added to the A-site La³⁺ ions [9]. LSMO has a polycrystalline perovskite crystal structure with a rhombohedral crystal structure in which an X-ray diffraction peak is separated near the diffraction angle of 2θ=33° [10]. However, LSBMO showed an orthorhombic crystal structure with a single peak, which is believed to be due to decrease of the distortion of the unit lattice caused by the addition of the smaller ionic radius Bi³⁺ (0.117 nm) [11].

Fig. 2 shows the surface (a–e) and cross-sectional (f) microstructures of the LSBMO films according to the amount of Bi³⁺. The surface SEM images show almost homogeneous and uniform microstructures. The average grain size tended to increase as the amount of Bi³⁺

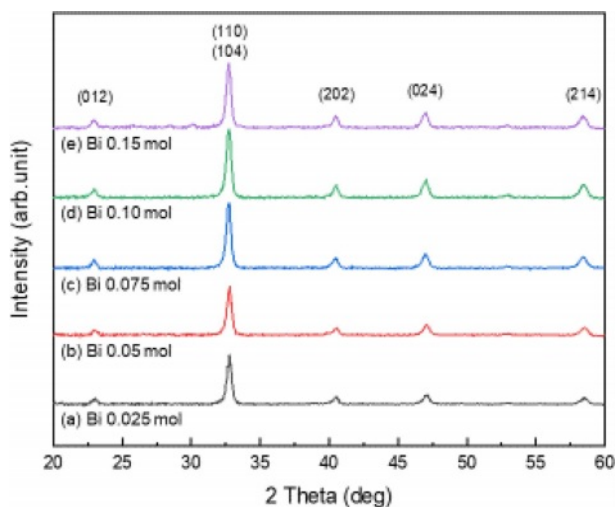


Fig. 1. X-ray diffraction patterns of LSBMO films with variation of Bi³⁺ contents.

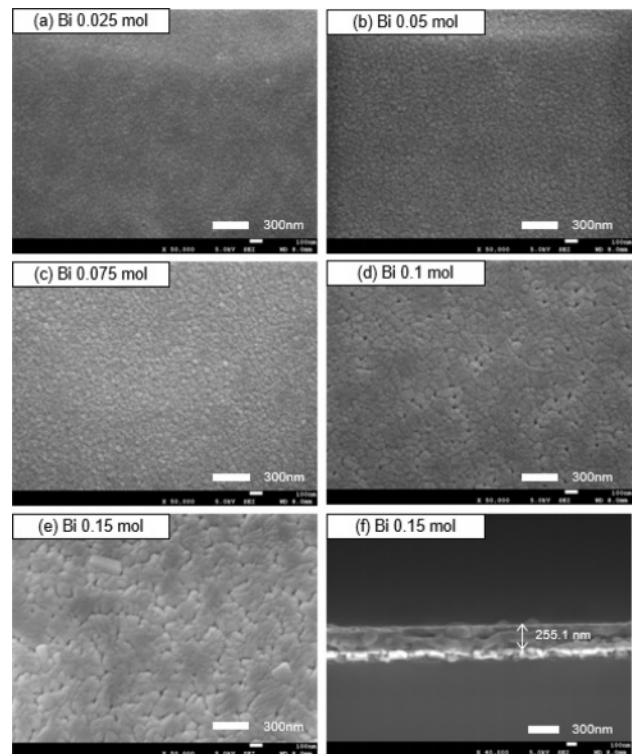


Fig. 2. (a)–(e) surface and (f) cross-sectional SEM micrographs of LSBMO films with variation of Bi³⁺ contents.

increased, and some small surface pores were observed in the La_{0.7}Sr_{0.15}Bi_{0.15}MnO₃ film. This is believed to be due to the volatilization of Bi (m.p.=825 °C), which has a low melting point, or surface irregularities due to the growth of grains in the vertical direction of the substrate [12]. The amount of Bi³⁺ did not seem to influence the thickness of the LSBMO films, and the average thickness for all films was approximately 230 to 260 nm.

Fig. 3 shows the XPS analysis results of the Mn 2p_{3/2} orbit according to the Bi³⁺ amount in the LSBMO films. The binding energies of Mn²⁺, Mn³⁺ and Mn⁴⁺ ions were not observed as dependent on the amount of Bi³⁺, and were approximately 640.3 to 641.0 eV, 641.5 to 642.16 eV and 644.23 to 644.6 eV, respectively. However, as the Bi³⁺ amount increased, the Mn⁴⁺/Mn³⁺ ratio tended to decrease, which is believed to be because the generation of Mn⁴⁺ formed to maintain the electrical neutrality of the unit lattice was suppressed by Bi³⁺ addition [13]. The shape of the Mn 2p signal of the LSBMO films implies that the majority of the Mn ions are in the Mn³⁺ state.

Fig. 4 shows the resistivity of the LSBMO films according to the amount of Bi³⁺ and temperature. All films showed typical ferromagnetic metallic properties in which resistivity increases as temperature increases [14]. The resistivity at room temperature increased as the amount of Bi³⁺ increased, and decreased after reaching a maximum value of 0.023 mΩ·cm at 0.075 mol. In general, the electrical conduction properties of

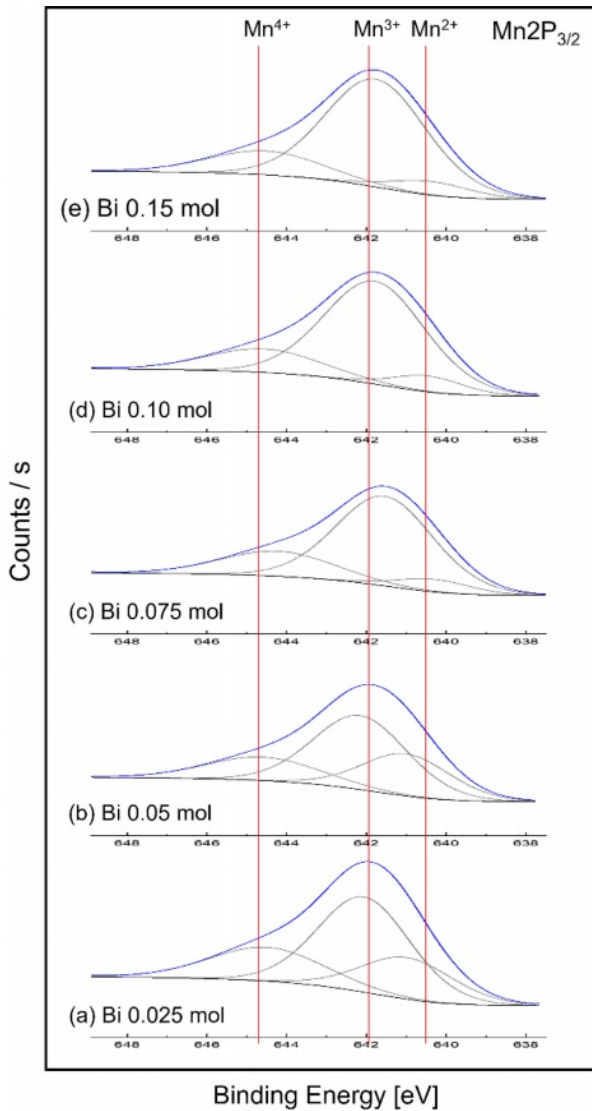


Fig. 3. XPS narrow scans of Mn 2P_{3/2} of LSBMO films with variation of Bi³⁺ contents.

LSMO result from the DE interaction of the outermost electrons of Mn³⁺ and Mn⁴⁺ via oxygen ions in a

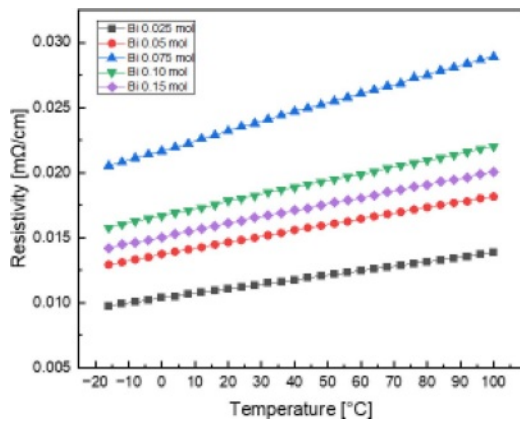


Fig. 4. Resistivity of LSBMO films with variation of Bi³⁺ contents.

perovskite lattice, and the electrical conduction properties are affected by the Mn-O bond distance and the Mn-O-Mn bond angle [15]. As shown in the XPS in Fig. 3, as the amount of Bi³⁺ increased, the resistivity increased because of the decrease in the Mn⁴⁺/Mn³⁺ ratio. However, when 0.10 mol or more Bi³⁺ is added, the Mn-O binding distance is decreased by the addition of Bi³⁺ (0.114 nm), which has a small ion radius, the hopping probability is increased and the resistivity is decreased. Further research is needed on the relationship between the transfer interaction of the e_g conduction electrons and the Jahn-Teller distortion caused by impurities with different valences and ionic radii to better understand the electrical conduction mechanism of LSMO materials.

Fig. 5 shows the temperature coefficient of resistance (TCR, $TCR = (1/RT)(dR_T/dT)$; R_T is resistance measured at T (°C)) and the B_{85/25}-value ($B_{35/25} = (\ln R_1 - \ln R_2)/(1/T_1 - 1/T_2)$, where R_1 and R_2 are the resistance measured at T_1 (25 °C) and T_2 (85 °C), respectively) according to the amount of Bi³⁺ in the LSBMO films. The TCR and B_{85/25}-value gradually increased as the amount of Bi³⁺ increased, and were the highest at 0.31%/°C and 299 K

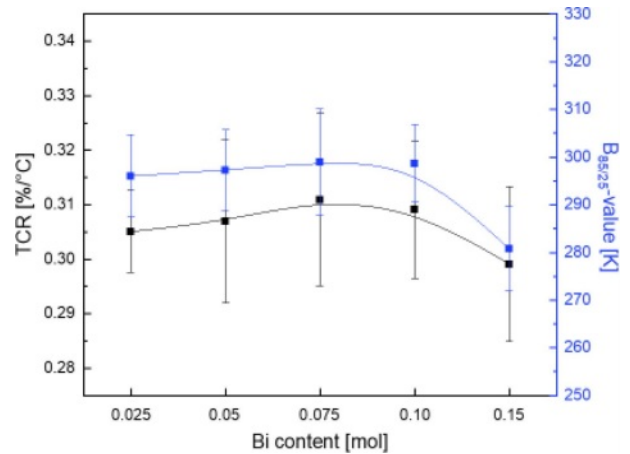
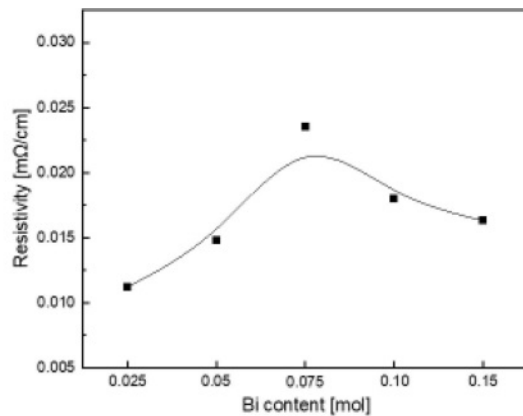


Fig. 5. TCR and B_{85/25}-value of LSBMO films with variation of Bi³⁺ contents.



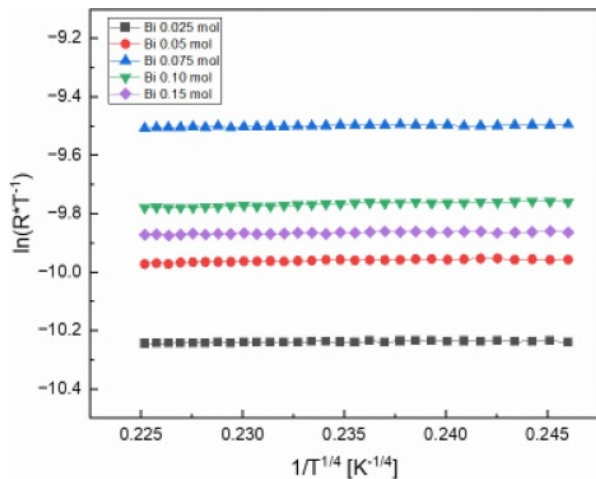


Fig. 6. $\ln(R \cdot T^{-1})$ vs $1/T^{1/4}$ curves of LSBMO films with variation of Bi³⁺ contents.

with 0.075 mol Bi³⁺, respectively. When 0.10 mol or more Bi³⁺ was added, the TCR and B_{85/25}-value tended to decrease. The crystal structure of perovskite R_{1-x}A_xMnO₃ is affected by the lattice distortion caused by the ionic radius of the cation located at the A-site, and its structural stability is expressed by the tolerance factor (*t*). An ideal cubic crystal structure without internal compressive or tensile strain is represented by *t* = 1, a rhombohedral structure is represented by 0.96 < *t* < 1 and an orthorhombic structure is represented by *t* < 0.96 [16]. The La_{0.7}Sr_{0.225}Bi_{0.075}MnO₃ film has a *t* of 0.9618, and is considered to have increased resistance change sensitivity due to thermal stimulation increases caused by the mixing of the rhombohedral and orthorhombic phases, resulting in an excellent TCR and B-value.

When the amount of Bi³⁺ added was 0.10 mol or more, the electrical conduction properties decreased due to the scattering effect of carriers in the pores and grain boundaries, as observed in Fig. 2 [17].

The electrical conduction of LSMO materials depends on the Mn⁴⁺/Mn³⁺ ratio and the Mn-O-Mn hopping probability. Polaron hopping transport according to temperature changes is generally expressed as follows [18]: $R = C_o T_o \exp(T_o/T)^p$, where *R* is the resistance, *C_o* is the constant, *T* is the temperature, and *T_o* is the characteristic temperature. $\alpha = p = 1$ is the nearest-neighbor hopping (NNH) model, and $\alpha = 4p$ is the variable-range hopping (VRH) model. The VRH model describes the jump of charge carriers from a localized state to another state [13]. Fig. 6 shows the $\ln(R \cdot T^{-1})$ vs $1/T^{1/4}$ curve of the LSBMO films according to the amount of Bi³⁺, and shows good linearity in all films. This is thought to be due to electron-electron scattering and electron-phonon interactions caused by thermal stimulation in the polycrystalline thin films in which crystal grains, grain boundaries and point defects are distributed [19, 20].

Conclusion

La_{0.7}Sr_{0.3-x}Bi_xMnO₃ (0.025 ≤ *x* ≤ 0.15) films showed the typical polycrystalline perovskite XRD patterns without any observable preferential orientations or impurities. The Mn⁴⁺/Mn³⁺ ratio tended to decrease as the amount of Bi³⁺ added increased. The shape of the Mn 2p orbital XPS signal of LSBMO films indicated that the majority of the Mn ions were in the Mn³⁺ state. The LSBMO films showed ferromagnetic metallic characteristics, and the electrical resistance properties are considered to be mutually affected by the transfer interaction of the e_g conduction electrons and Jahn-Teller distortion of the unit lattice. The electrical conduction mechanism of the LSBMO films showed VRH model dependence due to electron-electron scattering and electron-phonon interactions via thermal stimulation in the polycrystalline films in which crystal grains, grain boundaries, and point defects are distributed. La_{0.7}Sr_{0.225}Bi_{0.075}MnO₃ film exhibited a good TCR and B_{85/25}-value, and is expected to be applicable for thermistor devices.

Acknowledgements

This research was supported by Basic Science Research Program through the National Research Foundation of Korea (NRF) funded by the Ministry of Education (2020R1A6A1A03038697). This work was supported by the National Research Foundation of Korea (NRF) grant funded by the Korea government (MSIT) (2021R111A3052426), and the Technology Innovation Program (20020478) funded By the Ministry of Trade, Industry & Energy (MOTIE, Korea).

References

1. C. Sen, G. Alvarez, and E. Dagotto, Phys. Rev. Lett. 98 (2007) 127202.
2. J. Hemberger, A. Krimmel, T. Kurz, H. von Nidda, V. Inanov, A. Mukhin, A. Balbashov, and A. Loidl, Phys. Rev. B 66 (2002) 094410.
3. H. Heremans, J. Phys. D: Appl. Phys. 26[8] (1993) 1149-1168.
4. L. Mechin, J.M. Routoure, B. Guillet, F. Yang, S. Flament, D. Robbes, and R.A. Chakalov, Appl. Phys. Lett. 87 (2005) 204103.
5. A. Urushibara, Y. Moritomo, T. Arima, A. Asamitsu, G. Kido, and Y. Tokura, Phys. Rev. B 51[20] (1995) 14103-14109.
6. L.M. Rodriguez-Martinez and J.P. Attfield, Phys. Rev. B 58[5] (1998) 2426-2429.
7. X. Liu, Z. Jiao, K. Nakamura, T. Hatano, and Y. Zeng, J. Appl. Phys. 87[5] (2000) 2431-2436.
8. S. Kumari, N. Mottaghi, C.Y. Huang, R. Trappen, G. Bhandari, S. Yousefi, G. Cabrera, M.S. Seehra, and M.B. Holcomb, Sci. Rep. 10 (2020) 3659.
9. P. Norby, I.G. Anderson, E. Anderson, and N. Anderson, J. Solid State Chem. 119[1] (1995) 191-196.

10. A. Elghoul, A. Krichene, N.C. Boudjada, and W. Boujelben, *Ceramics International* 44[11] (2018) 12723-12730.
11. N. Nedelko, S. Lewinska, A. Pashchenko, I. Radelytskyi, R. Diduszko, E. Zubov, W. Lisowski, J.W. Sobczak, K. Dyakonov, A. S'lawska-Waniewska, V. Dyakonov, and H. Szymczak, *J. Alloys Compd.* 640 (2015) 433-439.
12. S. Wu, X. Wei, X. Wang, H. Yang, and S. Gao, *J. Mater. Technol.* 26[5] (2010) 472-476.
13. H. Xie, H. Huang, N. Cao, C. Zhou, D. Niu, and Y. Gao, *Phys. B: Cond. Matt.* 477[15] (2015) 14-19.
14. A. Elghoul, A. Krichene, and W. Boujelben, *J. Phys. Chem. Solid* 98 (2016) 263-270.
15. L.M. Rodriguez-Martinez and J.P. Attfield, *Phys. Rev. B* 54[22] (1996) 622- 625.
16. J.P. Zhou, J.T. McDevitt, J.S. Zhou, H.Q. Yin, J.B. Goodenough, Y. Gim, and Q.X. Jia, *Appl. Phys. Lett.* 75[8] (1999) 1146-1148.
17. X. Yu, S. Jin, H. Li, X. Guan, X. Gu, and X. Liu, *Appl. Surf. Sci.* 570 (2021) 151221.
18. N. Mott, *Conduction in Non-Crystalline Materials*, Clarendon, Oxford, 1993.
19. D.C. Worledge, G. Jeffrey Snyder, M.R. Beasley, T.H. Geballe, Ron Hiskes, Steve DiCarolis, *J. Appl. Phys.* 80[9] (1996) 5158-5161.
20. M. Viret, L. Ranno, and J.M.D. Coey, *Phys. Rev. B* 55[13] (1997) 8067-8070.